Maxwell's equations

459

VIII. A Dynamical Theory of the Electromagnetic Field. By J. CLERK MAXWELL, F.R.S.

Received October 27,—Read December 8, 1864.

PART I.—INTRODUCTORY.

(1) The most obvious mechanical phenomenon in electrical and magnetical experiments is the mutual action by which bodies in certain states set each other in motion while still at a sensible distance from each other. The first step, therefore, in reducing these phenomena into scientific form, is to ascertain the magnitude and direction of the force acting between the bodies, and when it is found that this force depends in a certain way upon the relative position of the bodies and on their electric or magnetic condition, it seems at first sight natural to explain the facts by assuming the existence of something either at rest or in motion in each body, constituting its electric or magnetic state, thing either at a distance according to mathematical laws.

Ŧ	or	Electromagn	netic Mo	menti	ım								• 1		F	G	H	
	22	Magnetic Ir	ntensity												α	β	7	
	"	Electromoti	ve Force												P	Q.	R	
	"	Current due	to true	condu	acti	on.									p	q	2	
	,,	Electric Dis													1777	g	ħ	
	,,	Total Curre													1000		7	
	,,	Quantity of															-	
	,	Electric Pot																
ween these twenty quantities we have found twenty equations, viz.																		
		hree equatio								-	-					(B)		
		22		ectric														
		22		ectron														
		33		ectric												(E)		
		22	Ele	ectric	Res	ista	nce				•					1		
		22		tal Cu												(A)		
×	0	ne equation	of Free	Electr	icit	v .										(G)		
		23	Conti													1		
		35.5				- 17		ALTER I						200	(708)	()_		- 1

These equations are therefore sufficient to determine all the quantities which occur in them, provided we know the conditions of the problem. In many questions, however, only a few of the equations are required.

The equations

The variations of the electrical displacement must be added to the currents p, q, r to get the total motion of electricity, which we may call p', q', r', so that

$$p' = p + \frac{df}{dt},$$

$$q' = q + \frac{dg}{dt},$$

$$r' = r + \frac{dh}{dt},$$
(A)

Equations of Magnetic Force.

$$\mu lpha = rac{d \, \mathrm{H}}{d y} - rac{d \, \mathrm{G}}{d z},$$
 $\mu eta = rac{d \, \mathrm{F}}{d z} - rac{d \, \mathrm{H}}{d x},$
 $\mu \gamma = rac{d \, \mathrm{G}}{d x} - rac{d \, \mathrm{F}}{d y}.$

Similarly,
$$\frac{d\gamma}{dy} - \frac{d\beta}{dz} = 4\pi p'.$$

$$\frac{d\alpha}{dz} - \frac{d\gamma}{dx} = 4\pi q',$$

$$\frac{d\beta}{dx} - \frac{d\alpha}{dy} = 4\pi r'.$$
(C)

We may call these the Equations of Currents.

Equations of Electromotive Force.

$$P = \mu \left(\gamma \frac{dy}{dt} - \beta \frac{dz}{dt} \right) - \frac{dF}{dt} - \frac{d\Psi}{dx},$$

$$Q = \mu \left(\alpha \frac{dz}{dt} - \gamma \frac{dx}{dt} \right) - \frac{dG}{dt} - \frac{d\Psi}{dy},$$

$$R = \mu \left(\beta \frac{dx}{dt} - \alpha \frac{dy}{dt} \right) - \frac{dH}{dt} - \frac{d\Psi}{dz}.$$

$$(D)$$

Equations of Electric Elasticity,

$$\begin{array}{c}
P = kf, \\
Q = kg, \\
R = kh.
\end{array}$$
(E)

Equations of Electric Resistance,
$$\begin{array}{c}
P = - gp, \\
Q = - gq, \\
R = - gr.
\end{array}$$
(F)

Equation of Free Electricity,
$$\begin{array}{c}
e + \frac{df}{dx} + \frac{dg}{dy} + \frac{dh}{dz} = 0.
\end{array}$$
(G)

Equation of Continuity,
$$\frac{de}{dt} + \frac{dp}{dx} + \frac{dq}{dy} + \frac{dr}{dz} = 0.$$
(H)

Modern form of Maxwell's equations

(Heaviside 1884)

$$\nabla \cdot \underline{E} = \frac{\rho}{\varepsilon_0}$$

Gauss's law

$$\nabla \cdot \underline{B} = 0$$

No magnetic monopoles

$$\nabla \times \underline{B} = \mu_0 \underline{J} + \mu_0 \varepsilon_0 \frac{\partial \underline{E}}{\partial t}$$

Ampere's law

$$\nabla \times \underline{E} = -\frac{\partial \underline{B}}{\partial t}$$

Faraday's law